

Operadores hermitianos

- $\int_a^b \psi^*(x) \varphi(x) dx = (\psi, \varphi) \rightarrow$ notação p/ o produto escalar de duas funções no intervalo (a, b)

- $\int \psi^*(x) \left[\frac{d}{dx} \varphi(x) \right] dx = \left(\psi, \frac{d}{dx} \varphi \right)$

- valores médios :

$$\begin{cases} \langle x \rangle = \int_{-\infty}^{\infty} \psi^*(x) x \psi(x) dx = (\psi, x \psi) \\ \langle p_x \rangle = \int_{-\infty}^{\infty} \psi^*(x) \left[-i\hbar \frac{\partial \psi(x)}{\partial x} \right] dx = \left(\psi, -i\hbar \frac{\partial}{\partial x} \psi \right) = (\psi, p_x \psi) \\ \langle E \rangle = (\psi, H \psi) \end{cases}$$

operador adjunto

$$\begin{aligned} \left(\psi, \frac{d}{dx} \varphi \right) &= \int_{-\infty}^{\infty} \psi^* \frac{d\varphi}{dx} dx = \cancel{\psi^* \varphi} \Big|_{-\infty}^{\infty} - \int \varphi \frac{d\psi^*}{dx} dx = \int \left(-\frac{d\psi^*}{dx} \right) \varphi dx \\ \text{operador linear} &= \left(-\frac{d\psi}{dx}, \varphi \right) \implies \boxed{\left(\frac{d}{dx} \right)^\dagger = -\frac{d}{dx}} \quad (\text{operador adjunto}) \\ &\quad \downarrow \text{operador adjunto de } \frac{d}{dx} \end{aligned}$$

- Para q.q. operador linear A , existe o operador adjunto tal que:

$$(\psi, A\varphi) = (A^+\psi, \varphi) \Rightarrow \boxed{(A^+)^+ = A}$$

- Se $A^+ = A$, operador é dito hermitiano (auto adjunto)

$$(\psi, \underbrace{AB}_{\chi}\varphi) = ((AB)^+\psi, \varphi)$$

$$(A^+\psi, \chi) = (A^+\psi, B\varphi) = (B^+(A^+\psi), \varphi) = ((B^+A^+)\psi, \varphi)$$

$$\Rightarrow \boxed{(AB)^+ = B^+A^+}$$

$$(\psi, A\psi) = (A\psi, \psi)^* = (A^+\psi, \psi)$$

$$\text{Se } A^+ = A \text{ (hermitiano)} \Rightarrow (A\psi, \psi) = (A\psi, \psi)^* = (\psi, A\psi) \text{ (real)} \Rightarrow \underline{\text{valor médio real}}$$

$$(\psi, x\psi) = \int \psi^* (x\psi) dx = \int (x\psi^*) \psi dx = (x\psi, \psi) \Rightarrow \boxed{x^+ = x} \text{ (hermitiano)} \Rightarrow \langle x \rangle = \text{real}$$

$$\begin{aligned}
 - \quad (\Psi, p_x \Psi) &= \int_{-\infty}^{\infty} \Psi^* \left(-i\hbar \frac{\partial \Psi}{\partial x} \right) dx = -i\hbar \left\{ \cancel{\Psi^* \Psi} \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} \Psi \frac{\partial \Psi^*}{\partial x} dx \right\} = \int_{-\infty}^{\infty} \left(i\hbar \frac{\partial \Psi^*}{\partial x} \right) \Psi dx \\
 &= \int_{-\infty}^{\infty} \left(i\hbar \frac{\partial \Psi}{\partial x} \right)^* \Psi dx = \int_{-\infty}^{\infty} (p_x \Psi)^* \Psi dx = (p_x \Psi, \Psi) \Rightarrow \boxed{p_x^\dagger = p_x} \text{ (hermitiano)}
 \end{aligned}$$

$$- \quad (\Psi, p_x^2 \Psi) = (p_x^\dagger \Psi, p_x \Psi) = (p_x \Psi, p_x \Psi) = (p_x^\dagger p_x \Psi, \Psi) = (p_x^2 \Psi, \Psi) \geq 0 \Rightarrow p_x^{2\dagger} = p_x^2 \text{ (hermitiano)}$$

$$\Downarrow \\
 \boxed{H^\dagger = H} \text{ (hermitiano)}$$

$$\begin{aligned}
 - \quad A \Psi_n &= a_n \Psi_n \quad \begin{array}{l} \nearrow \text{autovalor} \\ \searrow \text{autofunção} \end{array} \Rightarrow \left\{ \begin{array}{l} (\Psi_n, A \Psi_n) = (A \Psi_n, \Psi_n) \text{ (hermitiano)} \\ (\Psi_n, a_n \Psi_n) = (a_n \Psi_n, \Psi_n) \\ a_n (\Psi_n, \Psi_n) = a_n^* (\Psi_n, \Psi_n) \Rightarrow a_n^* = a_n \text{ (autovalores reais)} \end{array} \right.
 \end{aligned}$$

operador hermitiano

$$(\Psi_l, A \Psi_n) - (A \Psi_l, \Psi_n) = (a_n - a_l) (\Psi_l, \Psi_n) = 0 \Rightarrow (\Psi_l, \Psi_n) = 0 \quad (l \neq n) \text{ (autofunções ortogonais)}$$

Equações de Ehrenfest

①

$$\frac{d\langle A \rangle}{dt} = \frac{d}{dt} (\Psi, A\Psi) = \frac{d}{dt} \int \Psi^* (A\Psi) dx$$

$$= \int \frac{\partial \Psi^*}{\partial t} (A\Psi) dx + \int \Psi^* \left(\frac{\partial A}{\partial t} \right) \Psi dx + \int \Psi^* A \frac{\partial \Psi}{\partial t} dx$$

$$= \frac{i}{\hbar} \int (H\Psi^*) (A\Psi) dx - \frac{i}{\hbar} \int \Psi^* A (H\Psi) dx + \left\langle \frac{\partial A}{\partial t} \right\rangle$$

$\Downarrow$ hermitian

$$= \frac{i}{\hbar} \int \Psi^* H A \Psi dx - \frac{i}{\hbar} \int \Psi^* A H \Psi dx + \left\langle \frac{\partial A}{\partial t} \right\rangle$$

$$= \frac{i}{\hbar} \int \Psi^* \underbrace{(H A - A H)}_{[H, A]} \Psi dx + \left\langle \frac{\partial A}{\partial t} \right\rangle$$

$$i\hbar \frac{\partial \Psi}{\partial t} = H\Psi \Rightarrow \begin{cases} \frac{\partial \Psi}{\partial t} = -\frac{i}{\hbar} H\Psi \\ \frac{\partial \Psi^*}{\partial t} = \frac{i}{\hbar} H\Psi^* \end{cases}$$

$$\frac{d\langle A \rangle}{dt} = \frac{i}{\hbar} \langle [H, A] \rangle + \left\langle \frac{\partial A}{\partial t} \right\rangle$$

$$\frac{d\langle x \rangle}{dt} = \frac{i}{\hbar} \langle [H, x] \rangle$$

$$\boxed{\frac{d\langle x \rangle}{dt} = \frac{\langle p_x \rangle}{m}}$$

$$H = \frac{p_x^2}{2m} + V(x)$$

$$\Rightarrow [H, x] = \frac{1}{2m} [p_x^2, x]$$

$$= \frac{1}{2m} [p_x p_x, x] = \frac{1}{2m} [p_x p_x x - x p_x p_x + p_x x p_x - p_x x p_x + p_x x p_x]$$

$$= \frac{1}{2m} \left\{ p_x \underbrace{(p_x x - x p_x)}_{[p_x, x]} + \underbrace{(p_x x - x p_x)}_{[p_x, x]} p_x \right\}$$

$$[H, x] = \frac{1}{m} \underbrace{[p_x, x]}_{-i\hbar} p_x = -i\frac{\hbar}{m} p_x$$

$$\frac{d\langle p_x \rangle}{dt} = \frac{i}{\hbar} \langle [H, p_x] \rangle = - \left\langle \frac{\partial V}{\partial x} \right\rangle$$

$$\boxed{\frac{d\langle p_x \rangle}{dt} = \left\langle -\frac{\partial V}{\partial x} \right\rangle}$$

$\Downarrow$

$$\boxed{m \frac{d^2 \langle x \rangle}{dt^2} = \left\langle -\frac{\partial V}{\partial x} \right\rangle}$$

$$[H, p_x] = [V(x), p_x]$$

$$V(x)(p_x \psi) = V(x) \left(-i\hbar \frac{\partial}{\partial x} \psi \right) = -i\hbar V(x) \frac{\partial \psi}{\partial x}$$

$$p_x (V(x) \psi) = -i\hbar \frac{\partial}{\partial x} (V \psi) = -i\hbar \left(\frac{\partial V}{\partial x} \right) \psi + i\hbar V \frac{\partial \psi}{\partial x}$$

$$V(x) p_x \psi - p_x V(x) \psi = i\hbar \left(\frac{\partial V}{\partial x} \right) \psi$$

$$[H, p_x] = [V(x), p_x] = i\hbar \frac{\partial V}{\partial x}$$

(3)

$$m \frac{d^2 \langle x \rangle}{dt^2} = \left\langle -\frac{\partial V}{\partial x} \right\rangle = \langle F(x) \rangle \neq F(\langle x \rangle) \quad F(x) = -\frac{\partial V}{\partial x}$$

$$\neq \text{2a lei}$$

Casos excepcionais : partícula livre e oscilador

$$F(x) = F(0) + F'(0)x + F''(0)\frac{x^2}{2} + F'''(0)\frac{x^3}{3!} + \dots$$

$$\begin{aligned} \langle F(x) \rangle &= F(0) + F'(0)\langle x \rangle + F''(0)\frac{\langle x^2 \rangle}{2} + F'''(0)\frac{\langle x^3 \rangle}{3!} + \dots \\ &= \left(-\frac{\partial V}{\partial x} \right)_0 - \left(\frac{\partial^2 V}{\partial x^2} \right)_0 \langle x \rangle + \frac{\partial^3 V}{\partial x^3} \dots \end{aligned}$$

$$\left\{ \begin{array}{l} \text{part. livre : } V(x) = c\bar{x} \Rightarrow \frac{\partial V}{\partial x} = 0, \frac{\partial^2 V}{\partial x^2} = 0, \dots \\ \text{osc. harmônico : } V(x) = \frac{1}{2}kx^2 \Rightarrow \left(\frac{\partial V}{\partial x} \right)_0 = 0, \left(\frac{\partial^2 V}{\partial x^2} \right)_0 = k, \left(\frac{\partial^3 V}{\partial x^3} \right)_0 = 0, \dots \end{array} \right. \Rightarrow \frac{d^2 \langle x \rangle}{dt^2} = 0 \Rightarrow \frac{d \langle x \rangle}{dt} = c\bar{t} \Rightarrow \langle p_x \rangle = c\bar{t}$$

$$\downarrow$$

$$\langle F(x) \rangle = -k \langle x \rangle = F \langle x \rangle$$

$$m \frac{d^2 \langle x \rangle}{dt^2} = -k \langle x \rangle$$